

NTIC FILE COPY

AD-A224 830

NAME OF CONTRACTOR:

Paul M. Racciah

CONTRACT NUMBER:

DAAK 70-83-K-0047

EFFECTIVE DATE OF CONTRACT:

2/1/83

EXPIRATION DATE OF CONTRACT:

1/31/86

REPORTING PERIOD: */Title*

FIFTH QUARTER *for DAAK 70-83-K 0047*

PRINCIPAL INVESTIGATOR:

PAUL M. RACCAH

PHONE NUMBER:

(312) 996-3403

DTIC
ELECTE
AUG 06 1990
S D

THE VIEWS AND CONCLUSIONS CONTAINED IN THIS DOCUMENT ARE THOSE OF THE
AUTHORS AND SHOULD NOT BE INTERPRETED AS NECESSARILY REPRESENTING THE OFFICIAL
POLICIES, EITHER EXPRESSED OR IMPLIED, OF THE U. S. GOVERNMENT.

DISTRIBUTION STATEMENT A
Approved for public release
Distribution Unlimited

90 08 03 006

FIFTH QUARTER

RESEARCH PLAN

This quarter has been devoted to the study of the physical significance of the parameter Γ which can be obtained by Electrolyte Electrorreflectance (EER) or by spectroscopic ellipsometry.

RESULTS

The generalized theory of electrorreflectance (which has been carried out under separate funding) has permitted us to reconcile the Γ obtained by electrorreflectance to the Γ obtained by ellipsometry.

However, it does remain that the value of Γ obtained either way can be as high as 120 meV as compared to the expected 70 meV. We have determined that this is due to the breakdown of the virtual crystal assumption and is in fact a quantitative measure of it. It turns out that the measured $\Gamma_m = \Gamma_0 + kT + \Delta\Gamma$ where Γ_m is the measured linewidth, Γ_0 is the expected 70 meV and $\Delta\Gamma$ measures the departure from the virtual crystal approximation. (See Attached Expose)



Accession For	
NTIS CRASH	<input checked="" type="checkbox"/>
DTIC TAB	<input type="checkbox"/>
Unannounced	<input type="checkbox"/>
Justification	
By <i>perform 50</i>	
Distribution	
Availability Codes	
Dist	Special
<i>A-1</i>	

SIMPLE APPROXIMATION FOR EFFECT OF ALLOYING ON THE PHENOMENOLOGICAL LINEWIDTH Γ

Consider that $E_{cv}(\vec{k})$ is not a unique, sharply defined energy, but that for the absorption of a photon "locally", it depends on the local concentration c_{loc} averaged over a cluster of N atoms on the Hg-Cd sublattice.

$$E_{cv}^{loc}(\vec{k}) \approx E_{cv}^{loc}(\vec{k}_{cr}) + \frac{\hbar^2}{2\mu} (\vec{k} - \vec{k}_{cr})^2;$$

i.e., the dominant source of variation in $E_{cv}^{loc}(\vec{k})$ is the variation in the local critical point energy

$$E_0(c_{loc}) \approx E_{cv}^{loc}(\vec{k}_{cr}).$$

Then, for $N \gg 1$,

$$P(E_0) \approx (\sigma\sqrt{2\pi})^{-1} \exp \{-[E_0 - E_0(c)]^2/2\sigma^2\}$$

$$\text{with } \sigma^2 = E_1^2 c(1-c)/N,$$

$$\text{where } E(c_{loc}) \approx E(c) + (c_{loc} - c) E_1.$$

This leads to a replacement of the lineshape

$$L(E, \vec{k}, \Gamma_0) = -[E - E_{cv}(\vec{k}) + i\Gamma_0]^{-1}$$

$$\text{by } L(E, \vec{k}, \Gamma_0) \approx \int_{-\infty}^{\infty} [E - E_{cv}(\vec{k}) - [E_0 - E_0(c)] + i\Gamma_0]^{-1} P(E_0) dE_0$$

The only simple analytic result is obtained by replacing the Gaussian probability $P(E_0)$ by a Lorentzian probability. If one does this and chooses the Lorentzian probability to have width

$$\Gamma^1 = \sqrt{2\sigma},$$

which follows from an expansion of e^{-u^2} as $[1 + u^2 + \dots]^{-1}$, one finds that Γ_0 is replaced by

$$\Gamma_m = \Gamma_0 + \sqrt{2\sigma} = \Gamma_0 + E_1 \sqrt{2c(1-c)/N} + kT$$

This gives the following table:

$\Gamma - (\Gamma_0 + kT)$.04eV	.06eV	.08eV	.10eV
N	200	88	50	32

A better numerical approximation leads to values of N approx 40% larger.